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Systematic study of nonlinear dynamics and dynamic aperture optimization

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Outline

- 1. Relevant SuperKEKB lattice parameters
- 2. Nonlinearities
- 4. Phase space plots
- 5. Nonlinear detuning
- 6. Dynamic aperture study
- 7. DA optimization
- 8. Conclusion

SuperKEKB LER Lattice



IR with final focus and vertical chromatic section. Bymax = 2.7 km

SuperKEKB LER relevant parameters

Betatron tunes	Q _x /Q _y	45.527/45.570
Compaction factor	α	2.75 10 ⁻⁴
Damping times (ms)	ĭ,∕ ĭ s	31.8/16.8
Horizontal emittance (nm-rad)	\mathcal{E}_{χ}	5.78
Energy spread	<i>σ</i> ₌∕Ε	8.2·10 ⁻⁴
Natural chromaticity	ξ χ0 / ξ y0	-107.4/-807.5
Corrected chromaticity	ξ χ / ξ _y	0.7/3.5
Betatron coupling	K	0.4%
IP betas (mm)	β ^ˆ _∕ ∕β ^ˆ _y	31.1/0.270
Beam size at IP (µm)	σ _* /σ _y	13.5/0.079

All plot data below are given for IP

Strong chromatic sextupoles (FF horizontal and vertical correction sections) + arc chromatic sextupoles. All sextupoles are arranged in pairs separated by –I

Fringe fields in quadrupoles (third order)

Fringe fields in bends (second order)

Kinematic terms (third order)

Nonlinearities - sextupoles



For kick sextupoles all high order aberrations are cancelled exactly



For thick sextupoles only the second order terms are cancelled but third and higher appear

$$x_{3} = -x_{0} - p_{x0}L - \frac{(K_{2}L)^{2}}{12} (x_{0}^{3} + x_{0}y_{0}^{2}) \cdot L^{2} + \dots$$

$$p_{x3} = -p_{x0} - \frac{(K_{2}L)^{2}}{6} (x_{0}^{3} + x_{0}y_{0}^{2}) \cdot L + \dots$$

$$y_{3} = -y_{0} - p_{y0}L - \frac{(K_{2}L)^{2}}{12} (x_{0}^{2}y_{0} + y_{0}^{3}) \cdot L^{2} + \dots$$

$$p_{y3} = -p_{y0} - \frac{(K_{2}L)^{2}}{6} (x_{0}^{2}y_{0} + y_{0}^{3}) \cdot L + \dots$$

Nonlinearities – quad fringe fields

Particle transport through the quadrupole fringe (leading term)

$$\begin{aligned} x &= x_0 + \frac{k_{10}}{12(1+\delta)} \left(x_0^3 + 3x_0 y_0^2 \right), \qquad p_x = p_{x0} - \frac{k_{10}}{4(1+\delta)} \left[p_{x0} \left(x_0^2 + y_0^2 \right) - 2p_{y0} x_0 y_0 \right], \\ y &= y_0 - \frac{k_{10}}{12(1+\delta)} \left(y_0^3 + 3x_0^2 y_0 \right), \qquad p_y = p_{y0} + \frac{k_{10}}{4(1+\delta)} \left[p_{y0} \left(x_0^2 + y_0^2 \right) - 2p_{x0} x_0 y_0 \right], \end{aligned}$$

For simulation we use more complex but symplectic expressions from SAD

Rough comparison of two effects:

$$\Delta x_{qe} = -\frac{k_1}{12} \cdot x_0^3 = -4 \cdot 10^{-5} (cm^{-2}) \cdot x_0^3 \qquad \text{FF QC1LPH}$$

$$\Delta x_{ts} = -\frac{(k_2 L^2)}{12} \cdot x_0^3 = -8 \cdot 10^{-6} (cm^{-2}) \cdot x_0^3 \qquad \text{Thick sextupole pair in vertical}$$

$$\text{IR chromatic section}$$

Quadrupole fringe field effect is stronger than the tick sextupole effect

Nonlinearities – kinematic term

As we expect, a bulk of the kinematic effects

$$H_k(s) = \frac{1}{8} \left(p_x^2 + p_y^2 \right)^2$$

should come from IP with $Beta_Y = 270$ um.

The kinematic effects are included in the simulation.

Phase space plot (horizontal)



Extremely small seed of the vertical motion reduces the horizontal aperture by factor 2 through the nonlinear coupling term

Phase space plot (vertical)



1D and 2D (with small seed of horizontal motion) vertical apertures have the same size. Vertical DA is defined by the vertical nonlinear terms

Tune amplitude dependence



Main effect is due to the quadrupole fringe fields!

S + fringe fields+ kinematic terms₁₁

Illustration: dynamic aperture for thin and thick sextupoles arranged in pairs separated by -I map. Quadrupole fringe fields are switched off.



Dynamic aperture - II

Dynamic aperture decrease due to the kinematic effects. **Off-energy** LER SuperKEKB $\Delta E/E = 0.5\%$ (RF ON) z, σ_z **90** 80 On-energy LER SuperKEKB 60 z, g 90 50 80 40 70 30 20 60 10 50 n 40 -20 20 n **Χ**, σ_x 30 LER SuperKEKB $\Delta E/E = 1.0\%$ (RF ON) 20 ь∾ Ň 10 80 0 -20 20 0 70 **Χ**, σ_x 60 50 40 30 ---- Without kinematic terms 20 ---- With kinematic terms 10 0 20 -20 0

Χ, σ_x



Dynamic aperture - IV

Conclusion before optimization

DA in rms beam size



(1) DA due to the quadrupole fringe field and thick sextupoles have the same order of magnitude
 (2) Among all the quadrupoles influence of OC1 is most crucial

(2) Among all the quadrupoles influence of QC1 is most crucial

(1) Correction of the thick sextupole effect (by weak sextupole correctors)

(2) Correction of the quadrupole fringe fields (how?)

One can try to cope with the quadrupole edges (third order aberration) by octupole correctors compensating relevant detuning terms

To improve vertical aperture we have to correct vertical detuning term

To improve horizontal aperture we have to correct coupling detuning term

Dynamic aperture optimization - II

Separate octupole correctors, detuning coefficients





(1) In IR, where the bulk of the perturbation sources are concentrated, By>>Bx

(2) In any case $\beta_x \cdot \beta_y \le \max(\beta_x^2, \beta_y^2)$ so we can not control effectively the coupling term

(3) But with the vertical one everything works well

Dynamic aperture optimization - III

Octupole correctors incorporated in the quadrupole



Transformation through this simple model shows that we can effectively compensate either direct or coupling terms for both planes. The octupole corrector strength for each octupole kick is

$$k_3 L_o = \pm k_1^2 L_q$$

Later we found that a single kick or the octupole corrector distributed along the quadrupole length work as well.

Dynamic aperture optimization - IV

Example of the local correction scheme



DA reduced by the edges of strong quadrupole

Octupole corrector open it either horizontally (~10 times) (coupling term correction) or vertically (~3 times)



As we have two knobs to control the horizontal (by local correction) and vertical (by global correction) aperture, the following procedure was applied

(1) Optimization of the horizontal aperture in main quadrupoles

(2) Optimization of the vertical aperture by two octupoles in the vertical section of the FF chromaticity correction

(3) Numerical optimization of all octupole correctors because in real lattice every corrector not only suppresses a required term but also affects others (though slightly)

Dynamic aperture optimization - VI

Optimization of the quadrupole fringe fields by octupole correctors

z, cm

Source	Nx	Ny
Fringe QC1	25	80
Fringe QC2	50	380
Sextupoles only	40	60
All	30	80
Corrected fringes	100	1300

Octupole correctors

Name	L, cm	$k_{3}L$, cm ⁻³	<i>B</i> ‴ , G/cm ³
OQC1LP	39	-8.56·10 ⁻⁶	-2.9
OQC1RP	28	–1.18·10 ⁻⁵	-5.6
OQC2LP	35	1.11·10 ⁻⁵	4.2
OQC2RP	35	1.11·10 ⁻⁵	4.2
OY1L=OY1R	10	-9.95·10 ⁻⁴	133
OY2L=OY2R	10	3.28·10 ⁻⁶	4.4



Black - initial aperture Blue – FF quadrupoles corrected Red – all quadrupoles corrected

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Dynamic aperture optimization - VII



Dynamic aperture optimization - VIII

Thick sextupole correction by two additional weak (strength = 3-7% of the main one) sextupole correctors (not applied yet to SuperKEKB)



(1) In SuperKEKB LER the dynamic aperture is limited in the first place by the fringe field in the strong FF quadrupoles

(2) In the second place by the finite length strong chromatic sextupoles

(3) The quadrupole fringe fields can be corrected by octupole correctors, individual and incorporated in the FF quadrupoles

(4) The next steps are: (a) including and optimization of the sextupoles, (b) global optimization of all effects for on- and off-momentum particle, (c) optical errors (!)