

Design of Damping Ring

Kikuchi, M., MAC2011, 7 Feb '11

Requirements on emittance and intensity

a. Injection aperture of LER

$$2J = 0.7 \mu\text{m} \quad (0.5 \mu\text{m, MAC10})$$

- Emittance of Injected beam

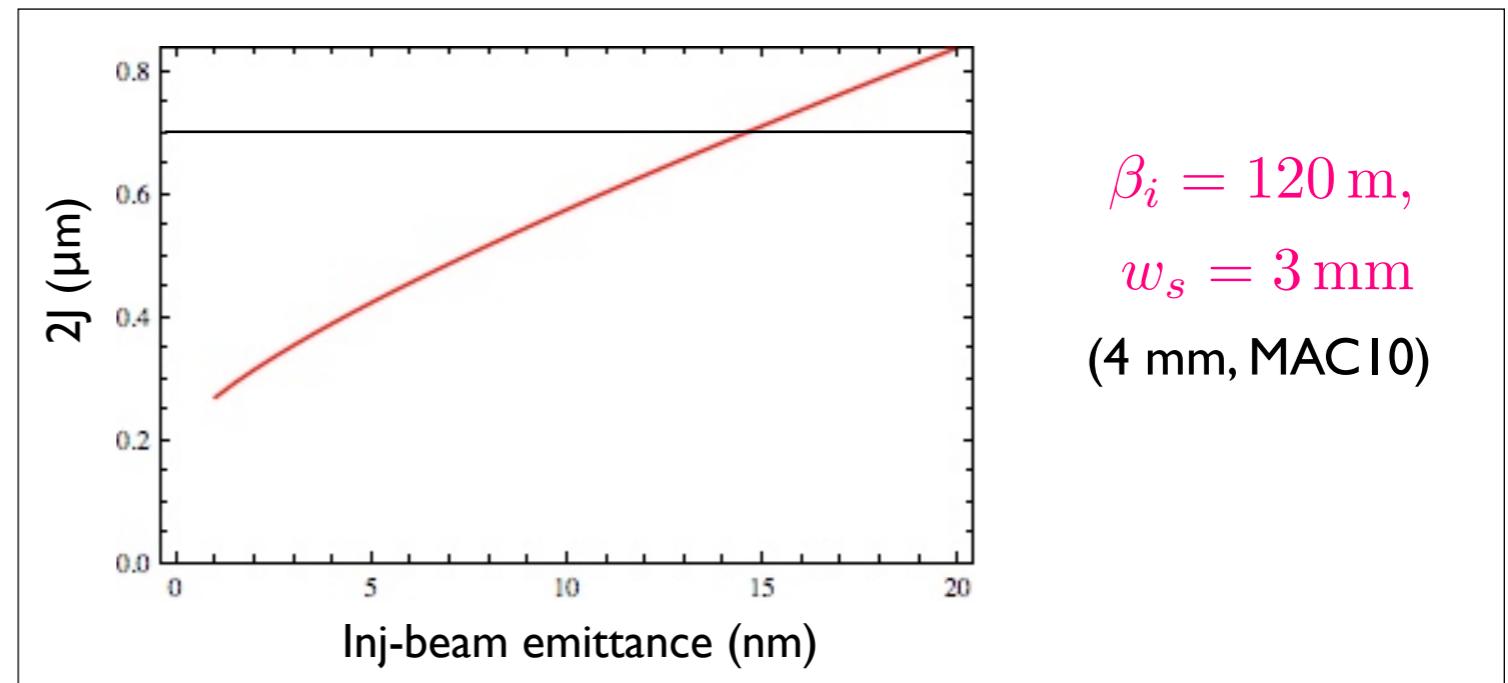
$$\epsilon_i \leq 14.5 \text{ nm} \quad (4 \text{ nm, MAC10})$$

b. Lifetime of LER

$$\tau \simeq 600 \text{ sec}$$

- Intensity of Injected beam

$$q \simeq 8 \text{ nC/pulse} = 4 \text{ nC/bunch}$$



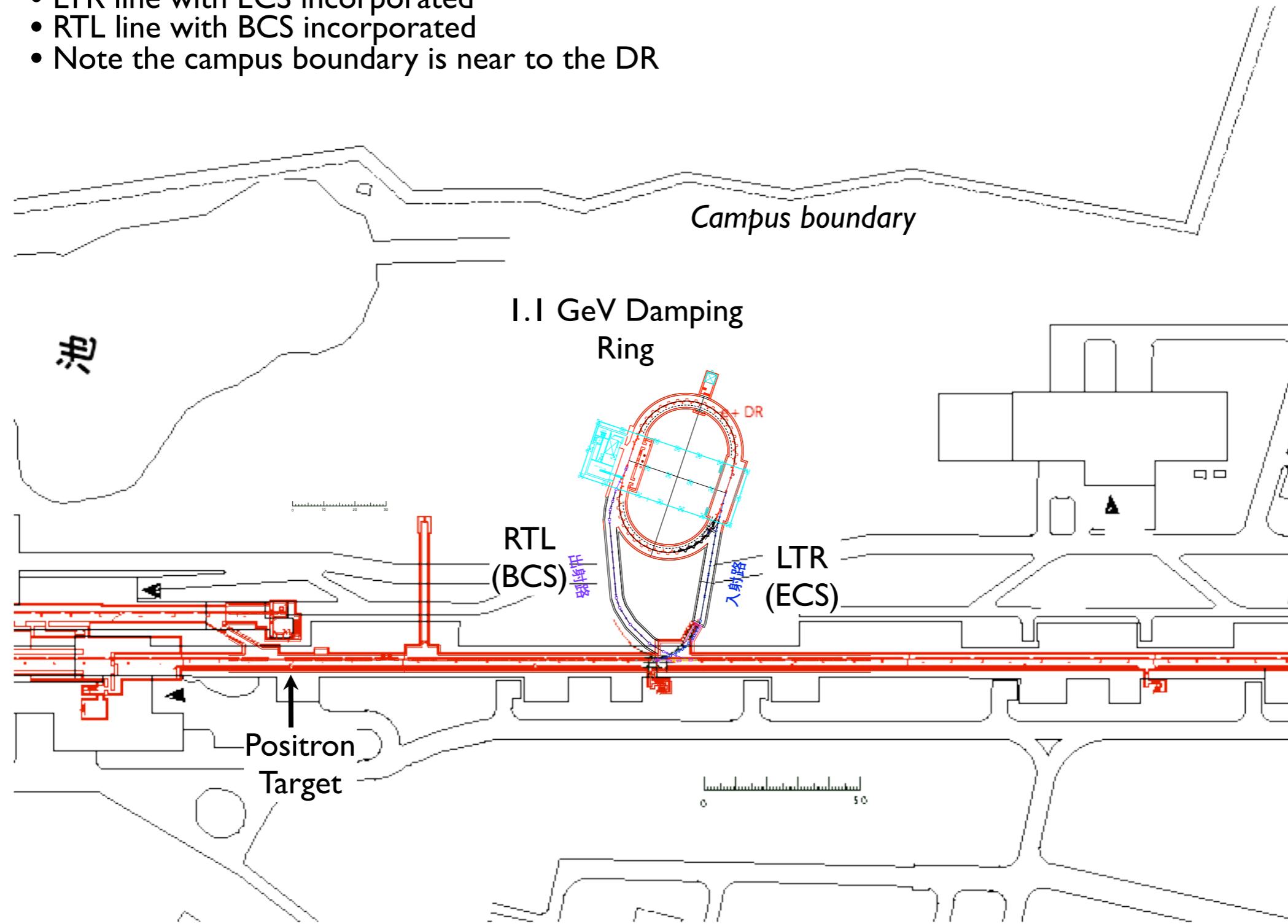
$$\dot{I} = I/\tau = 3.6 \text{ A}/600 \text{ sec} = 6 \text{ mA/sec}$$

$$\dot{Q} = \dot{I}/f_{rev} = 60 \text{ nC/sec} = 2.4 \text{ nC}@25 \text{ Hz}$$

This corresponds to 30 % injection efficiency

Layout of the System

- Positron target at sector I-4
- Capture + acceleration with L-band+S-band structure
- Extract from Linac at 1.1 GeV
- After 2 linac-pulse, re-inject to the Linac
- LTR line with ECS incorporated
- RTL line with BCS incorporated
- Note the campus boundary is near to the DR



Optics

Concept

- Large dynamic aperture
- Shorter damping time with lower bend field

FODO with Reverse Bend

$$\begin{aligned}\tau &= \frac{3T_0}{r_e \gamma^3 J_x I_2} = \frac{3}{2\pi c r_e J_x} \frac{\rho}{\gamma^3} C \frac{1-r}{1+|r|} \\ &= \frac{3}{2\pi c r_e J_x} \frac{\rho}{\gamma^3} \left(2\pi\rho + \frac{1-r}{1+|r|} L_1 \right)\end{aligned}$$

r : Bend ratio = $B1/B2$
(Normal FODO $\rightarrow r = -1$)

L_1 : Total length except bend length

L_1 is effectively reduced by a factor of
 $(1-r)/(1+|r|) = 0.48$

$r = 0.35, \rho = 2.7 \text{ m}, L_1 = 100 \text{ m}$

$$2\pi\rho = 17 \text{ m}$$

- Shorter damping time with lower bend field

$$B = 1.35 \text{ T}$$

Ring Parameters

Parameters of the Damping Ring

	1.1	GeV
No. of bunch trains/ bunches per train	2 / 2	
Circumference	135.5	m
Maximum stored current*	70.8	mA
Energy loss per turn	0.091	MV
Horizontal damping time	10.9	ms
Injected-beam emittance	1700	nm
Equilibrium emittance(h/v)	41.4 / 2.07	nm
Coupling	5	%
Emittance at extraction(h/v)	42.5 / 3.15	nm
Energy band-width of injected beam	± 1.5	%
Energy spread	0.055	%
Bunch length	6.5	mm
Momentum compaction factor	0.0141	
Number of normal cells	32	
Cavity voltage for 1.5 % bucket-height	1.4	MV
RF frequency	509	MHz
Inner diameter of chamber	32	mm
Bore diameter of magnets	44	mm

MAC10

1.0

- 8 nC/bunch (16 nC/pulse) is the ultimate goal
- The hardware design is based on this value.

12.7

2100

14 / 1.4

10

17.6 / 5.1

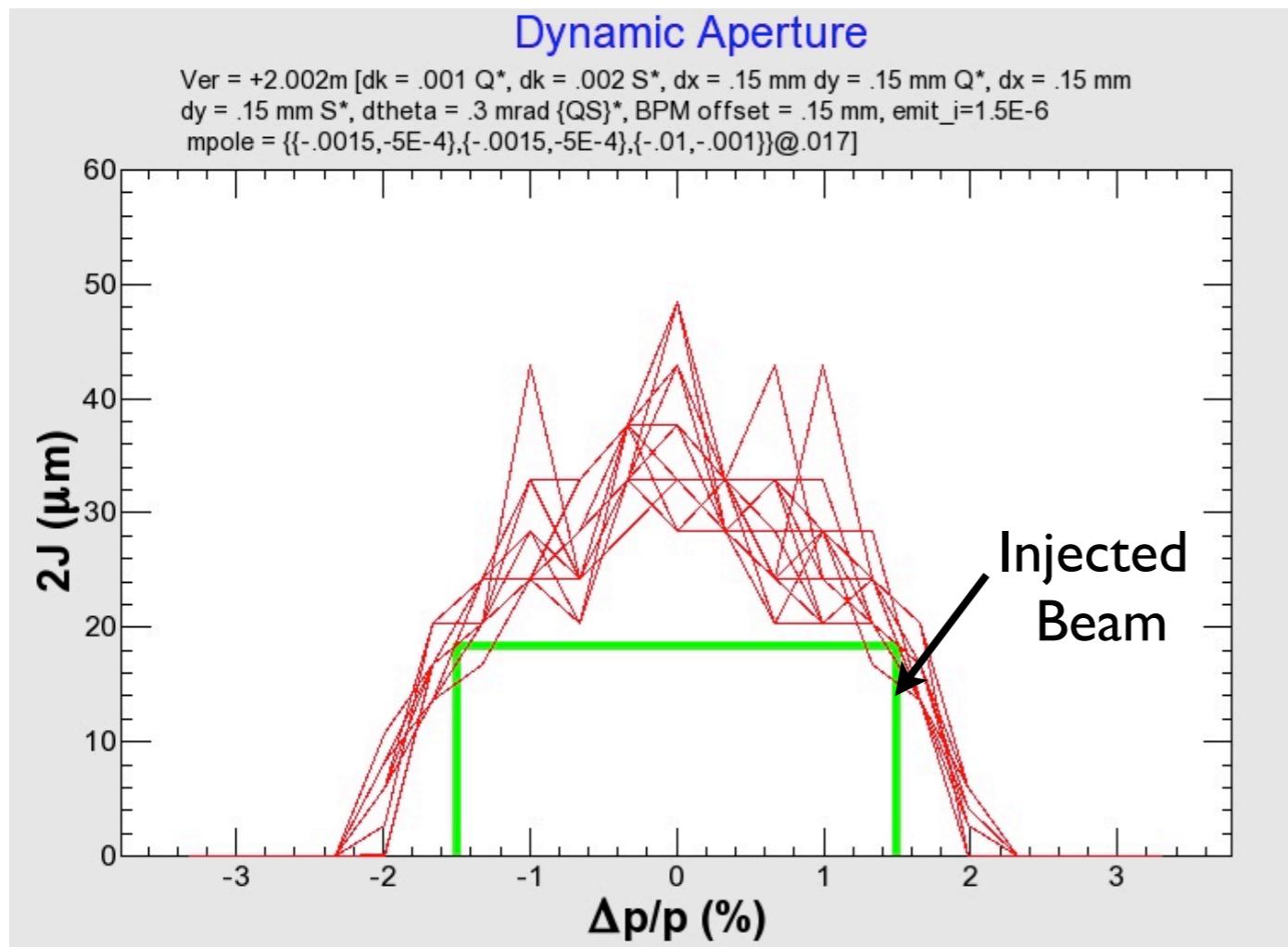
5.4

0.0019

0.26

* 8 nC/bunch

Dynamic aperture



Injected beam: $\epsilon = 1.5 \mu\text{m}$, 3.5σ
 $|\Delta p/p| = 1.5\%$

'Normal' errors

- Strength: 0.1 % for quads, 0.2 % for sexts
- Rotation: 0.3 mrad for quads and sexts
- Misalignment: 0.15 mm
- BPM offset: 0.15 mm
- Orbit correction

PLUS

Higher order multipoles
(Systematic, $\Delta B/B$ at $r=17 \text{ mm}$)

- Bend: K2 0.15 %, K4 0.05 %
 - Quad: K5 0.15 %, K9 0.05 %
 - Sext: K8 1.0 %, K14 0.1 %
- (Except quads in straight)

- Bends and quads are in fabrication based on these numbers

Electron Cloud Instability

- Theory

$$\rho_{e,th} = \frac{2 \ln 2\pi}{3\sqrt{2}} \frac{\gamma \nu_s \omega_e \sigma_z / c}{K Q r_e \beta L} \left(1 + \frac{\sigma_y}{\sigma_x}\right)$$

$$\omega_e^2 = \lambda_+ r_e c^2 / \sigma_y (\sigma_x + \sigma_y)$$

$$Q = \min(5, \omega_e \sigma_z / c) \quad K = 3$$

- Correction factor due to discrepancy between the theory and the simulation: 0.43 (simulation gives smaller threshold)

$$(\rho_e L)_{th} = 7.0 \times 10^{14} \text{ m}^{-2}$$

- Simulation results of photo-electron formation

Cloud density: ρ					
		Drift	Bend	Q + SX	unit
Length		73.2	36	26.8	m
$\delta_{MAX}=2$	SR=1	1.3	0.6	0.5	10^{12} m^{-3}
	SR=1	0.4	0.5	0.15	10^{12} m^{-3}
$\delta_{MAX}=1$	SR=0.1	0.15	0.11	0.03	10^{12} m^{-3}

- SR is photon-flux ratio to the design flux:
- SR=0.1 mimics the anti-chamber effect
- Electron Cloud is not proportional to the flux
- $\delta_{MAX}=1$ expected for TiN coated Al chamber

Electron Cloud Instability 2

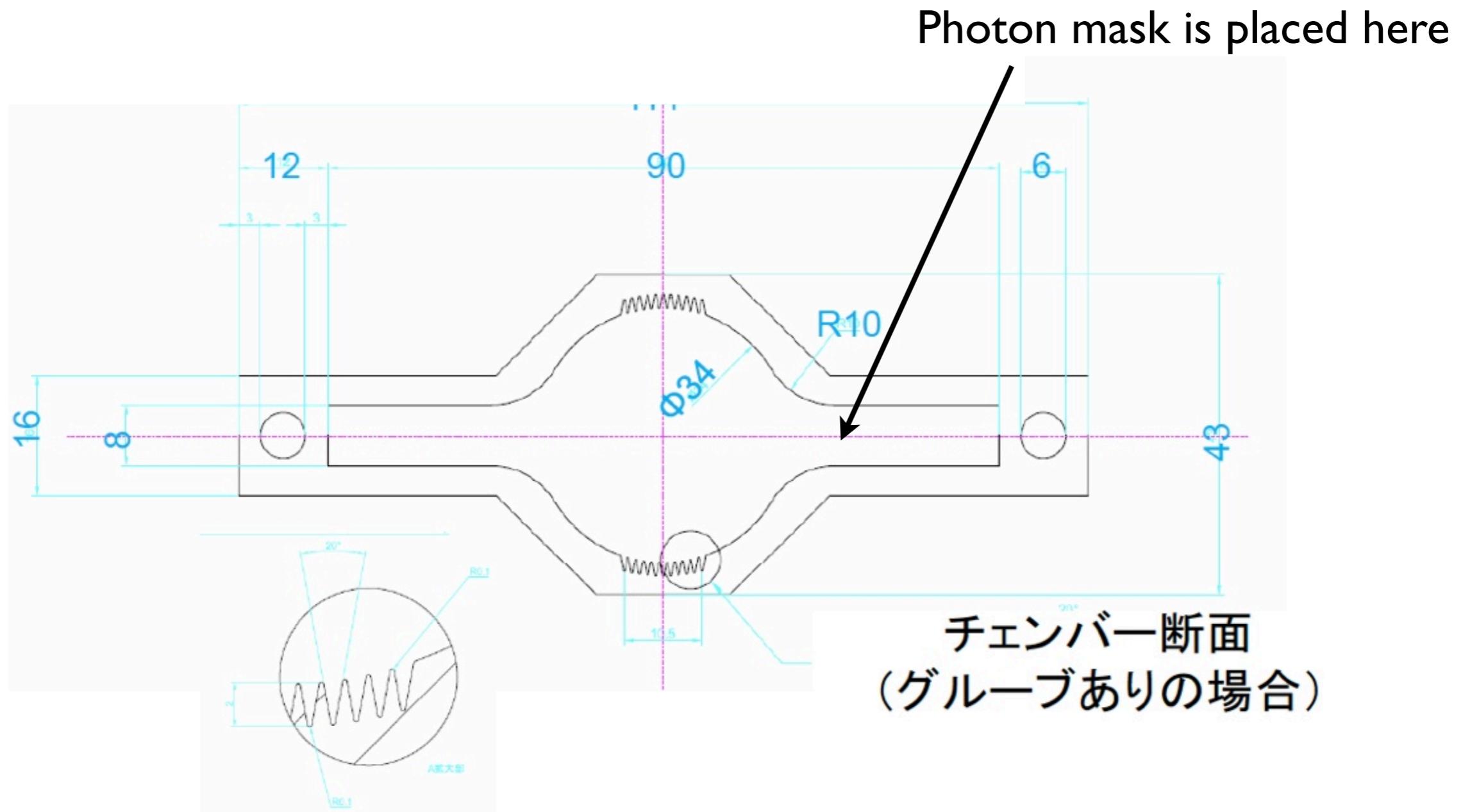
- Integrated electron density for $\delta_{max}=1$, SR=1

$$0.51 \times 10^{14} \text{ m}^{-2} \ll \rho_{e,th} L = 7.0 \times 10^{14} \text{ m}^{-2}$$

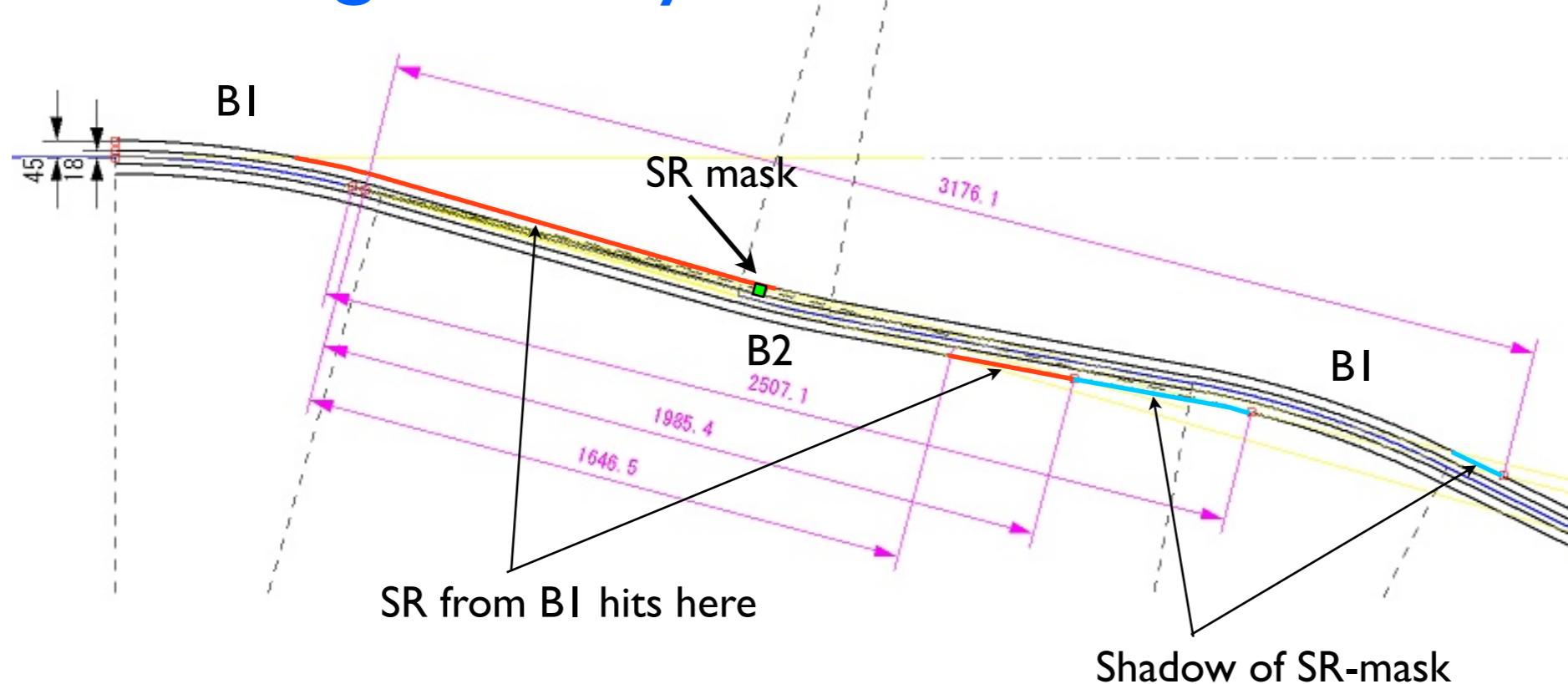
- Electron density is well below the threshold
- Electron cloud in reality may be larger than 2D simulation.
- Some mitigation technique will be incorporated in the vacuum design for insurance
- (Antechamber is needed for the impedance-free photon mask)

Chamber (basic plan)

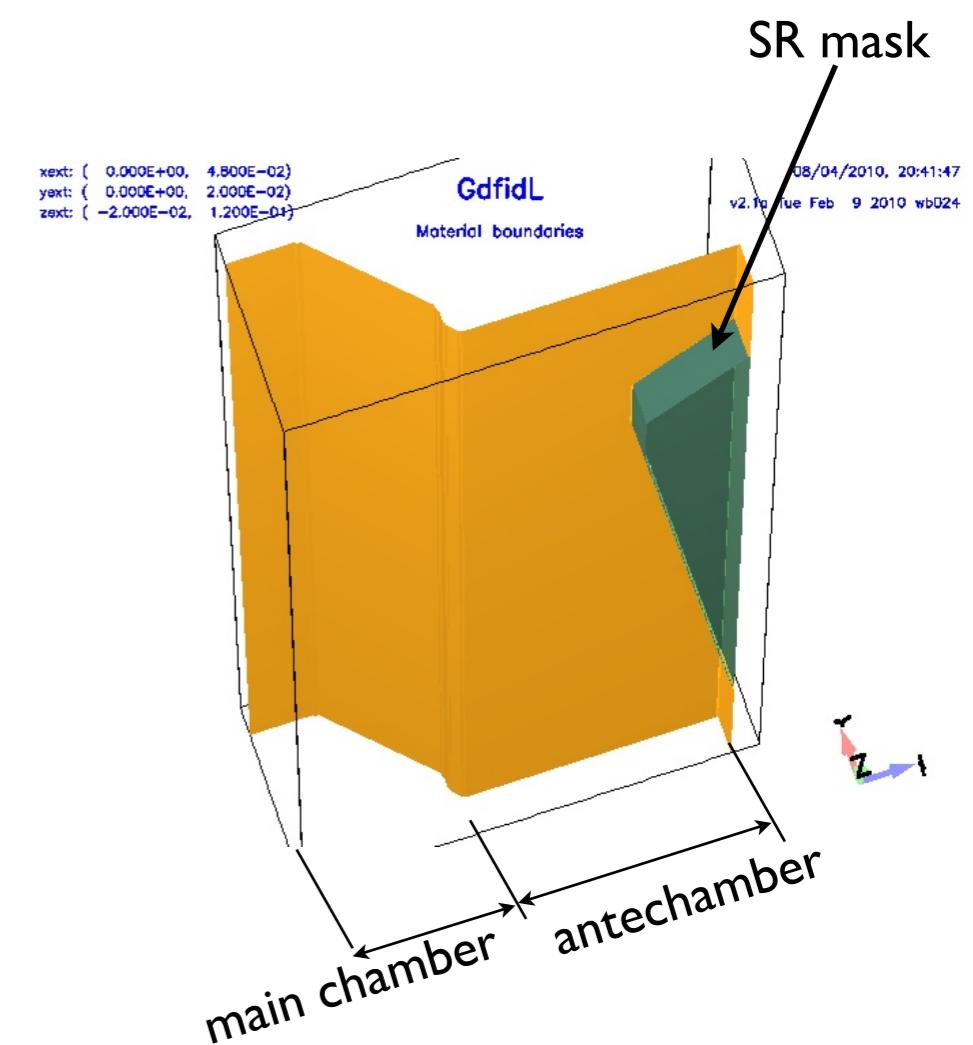
- Al alloy
- Antechamber with cooling-water channels
- Grooved surface on top and bottom of the bend chamber



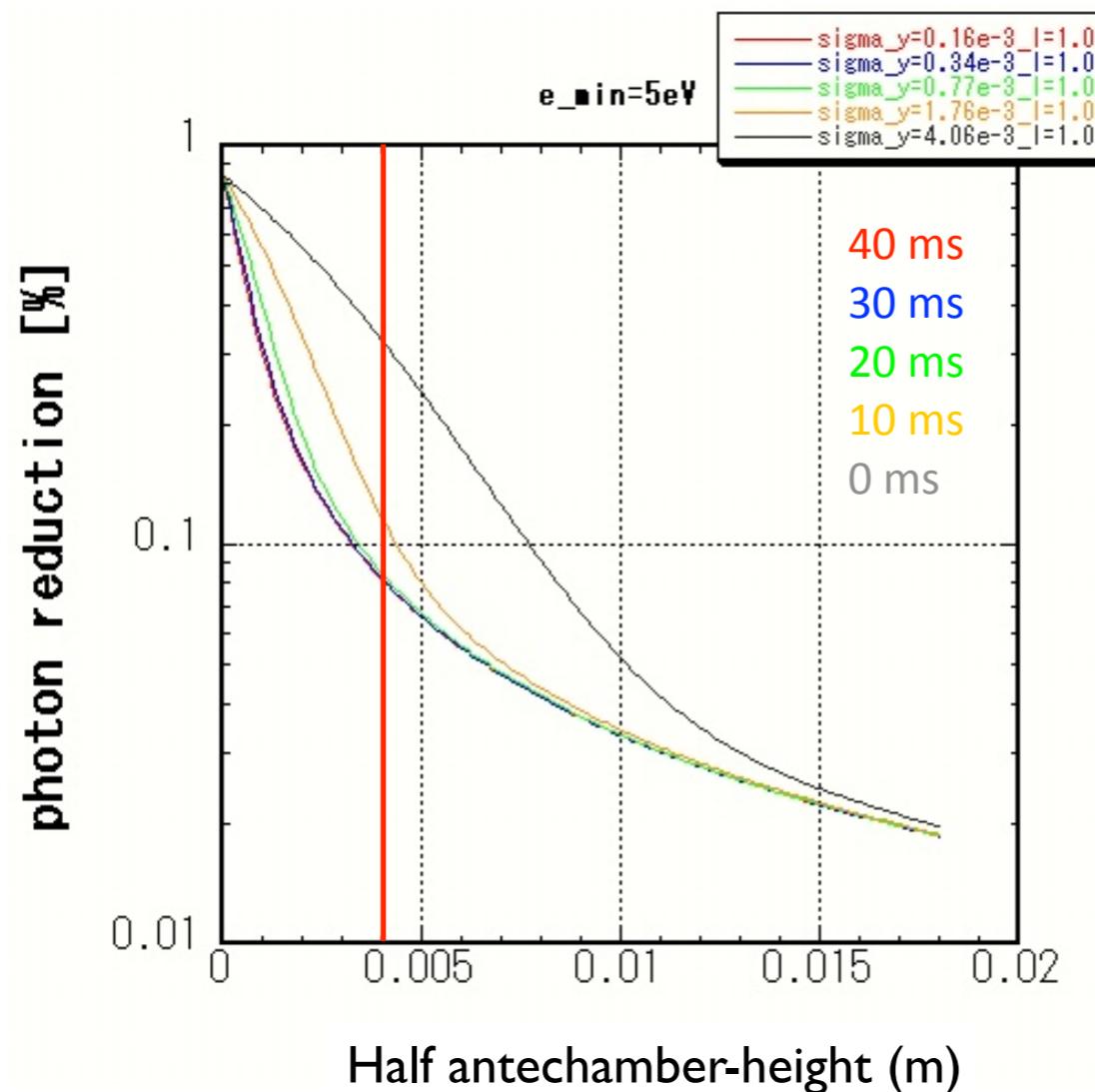
SR fan geometry for BI



- Antechamber + negative bend provides SR more chance to hit the downstream chamber
- 90 % of SR power is dropped within 1m from BI
- SR mask housed in the antechamber makes the shadow at more than 2 m from the source point
- The SR mask is useful also to prevent the SR from hitting bellows.
- The SR mask has negligibly small impedance; 1e-6 compared to the normal mask.



Acceptance of antechamber for the photons



- Vertical axis: Ratio of photons that hit surface of the main chamber
- Neglect photons $E < 5 \text{ eV}$
- Distance from source point is 1 m
- Antechamber height of 8 mm is sufficient to accept 90 % of photons in the antechamber if the store time is longer than 10 ms.

Transverse impedance

- Calculation results of the kick factor

2010.04.16

Component	No.		kick_x	N*kick_x	kick_y	N*kick_y
ARES	1		1.74E+13	1.74E+13	-1.70E+13	-1.70E+13
Bellows	88	X	2.93E+11	2.57E+13	-1.56E+08	
		Y	-1.51E+08		2.84E+11	2.50E+13
flange gap	176	X	4.73E+11	8.32E+13	1.06E+07	
		Y	-4.93E+07		5.06E+10	8.91E+12
Pumping port	176	X	1.98E+08	3.48E+10	-1.13E+08	1.99E+10
SR mask	176	XY	2.20E+06	3.87E+08	-2.63E+05	4.63E+07
Resistive wall	135.5	XY	-3.6E+12	4.88E+14	-3.6E+12	4.88E+14
BPM	82	X	0.571	46.82	0.483	39.606
Stripline kicker	1	X	33.7	33.7		
			V/Cm	6.14×10 ¹⁴	V/Cm	5.39×10 ¹⁴

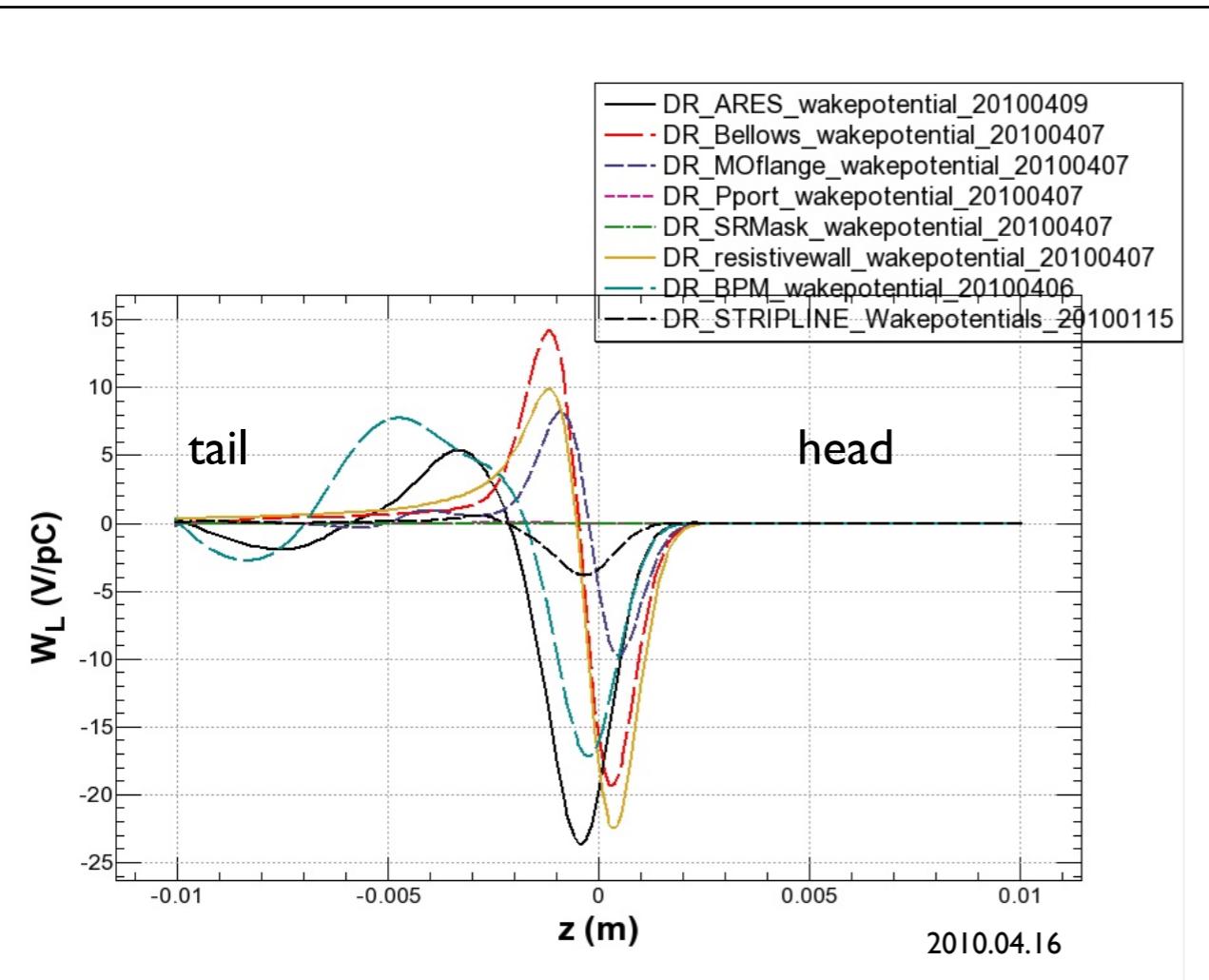
- Transverse microwave instability

$$I_{th} = \frac{8f_s E/e}{\sum \beta_i k_i} = 0.41 \text{ A; for x} \longrightarrow 183 \text{ nC/bunch}$$
$$0.31 \text{ A; for y} \longrightarrow 141 \text{ nC/bunch}$$

Longitudinal impedance

Vacuum chamber component

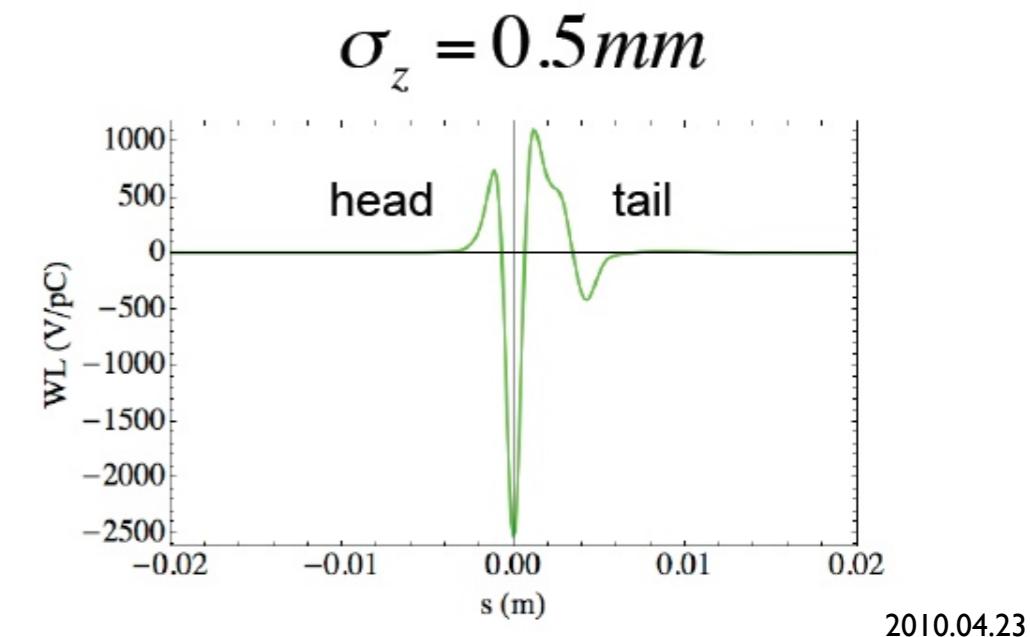
Shibata, K.



CSR

D. Zhou

- Stupakov's code
- Toroidal rectangular chamber (34×34)
+ infinite straight pipe
- Sum. of single bend wake

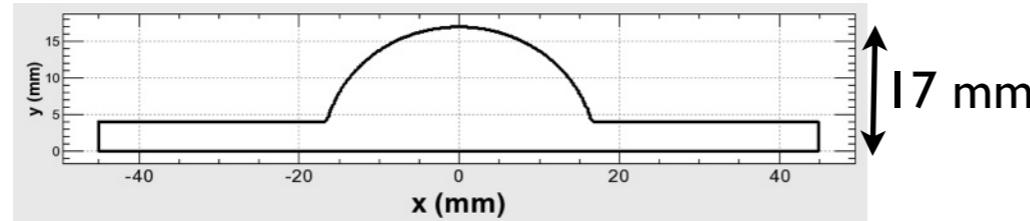


- CSR wake is 100 times larger than the ordinal vacuum-components wake!

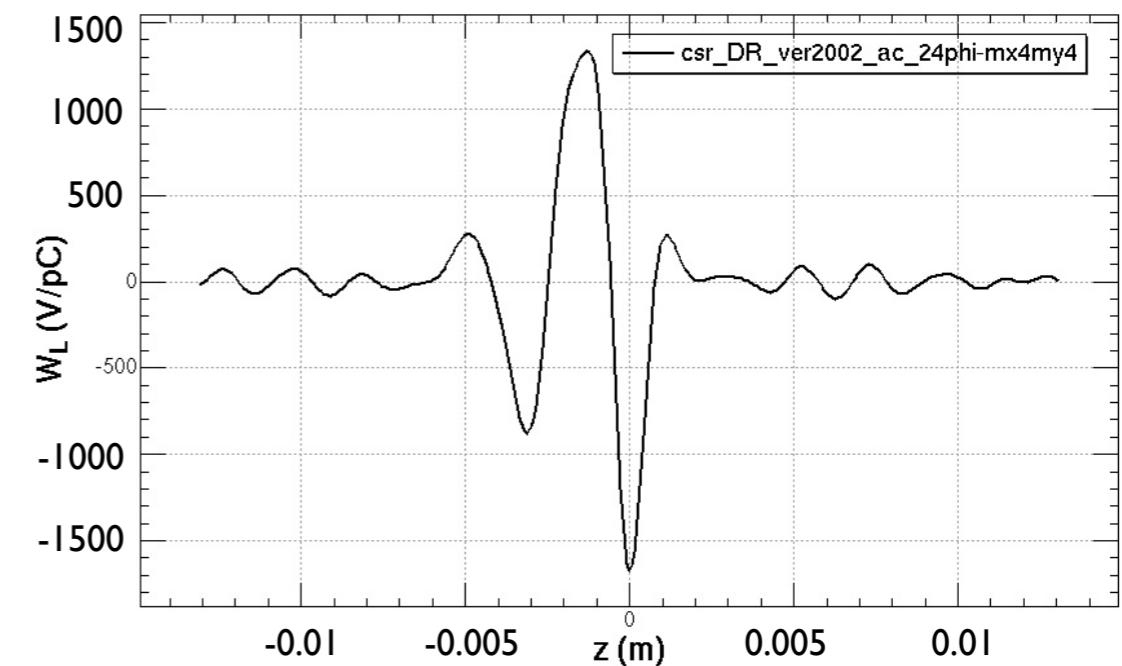
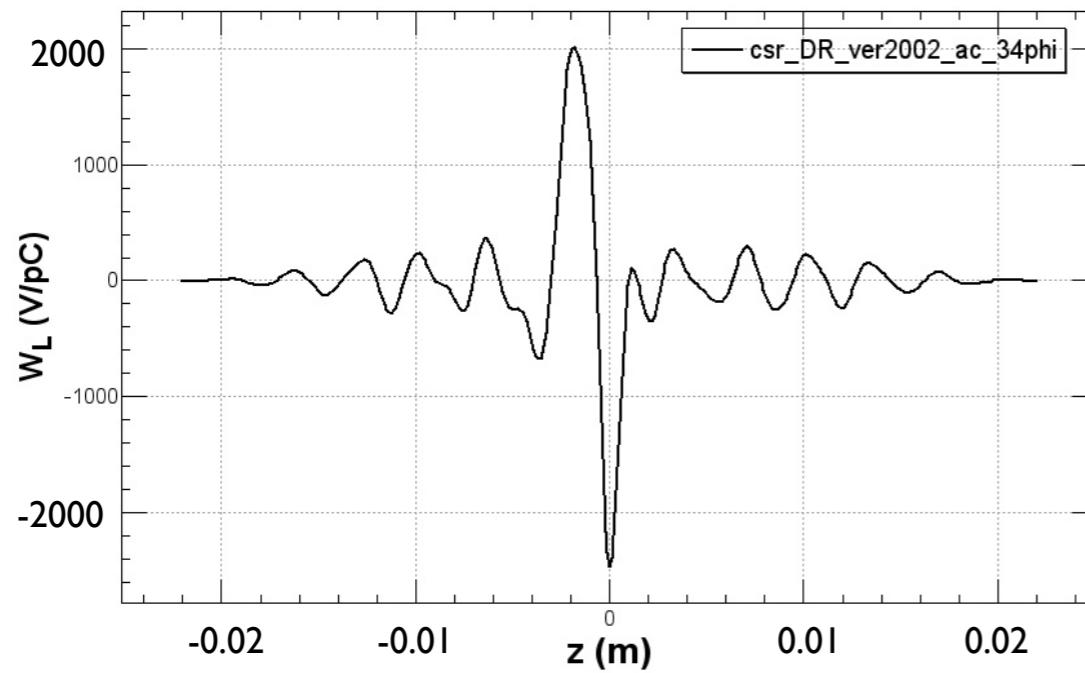
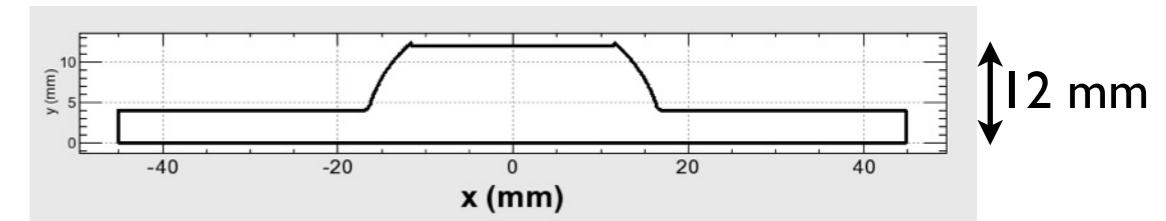
- CSR wake of proposed chambers

- Numerical Calculation - Oide's code
- Model : ante-chamber (single bend :ver2002)+ infinite straight pipe
 - summation over $32 \times B1 + 38 \times B2 + 6 \times B3 + 2 \times B4$
 - $\sigma z = 0.5\text{mm}$

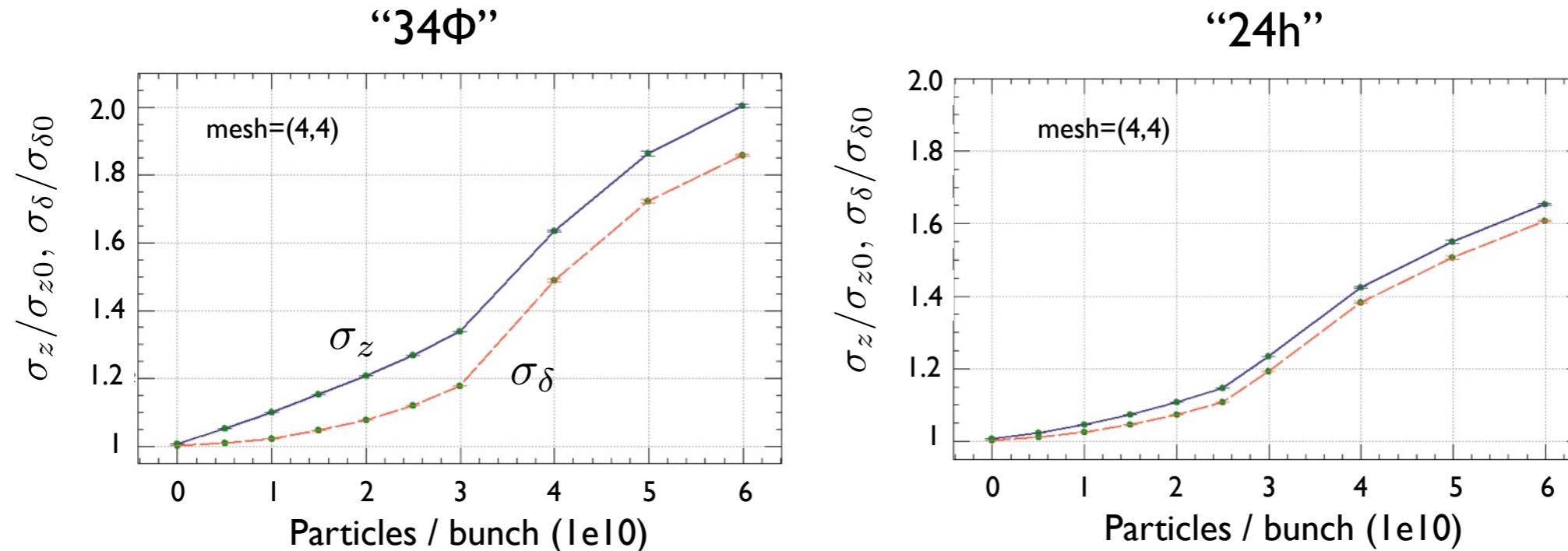
“34Φ”



“24h”



- Multi-particle longitudinal tracking using the CSR wake

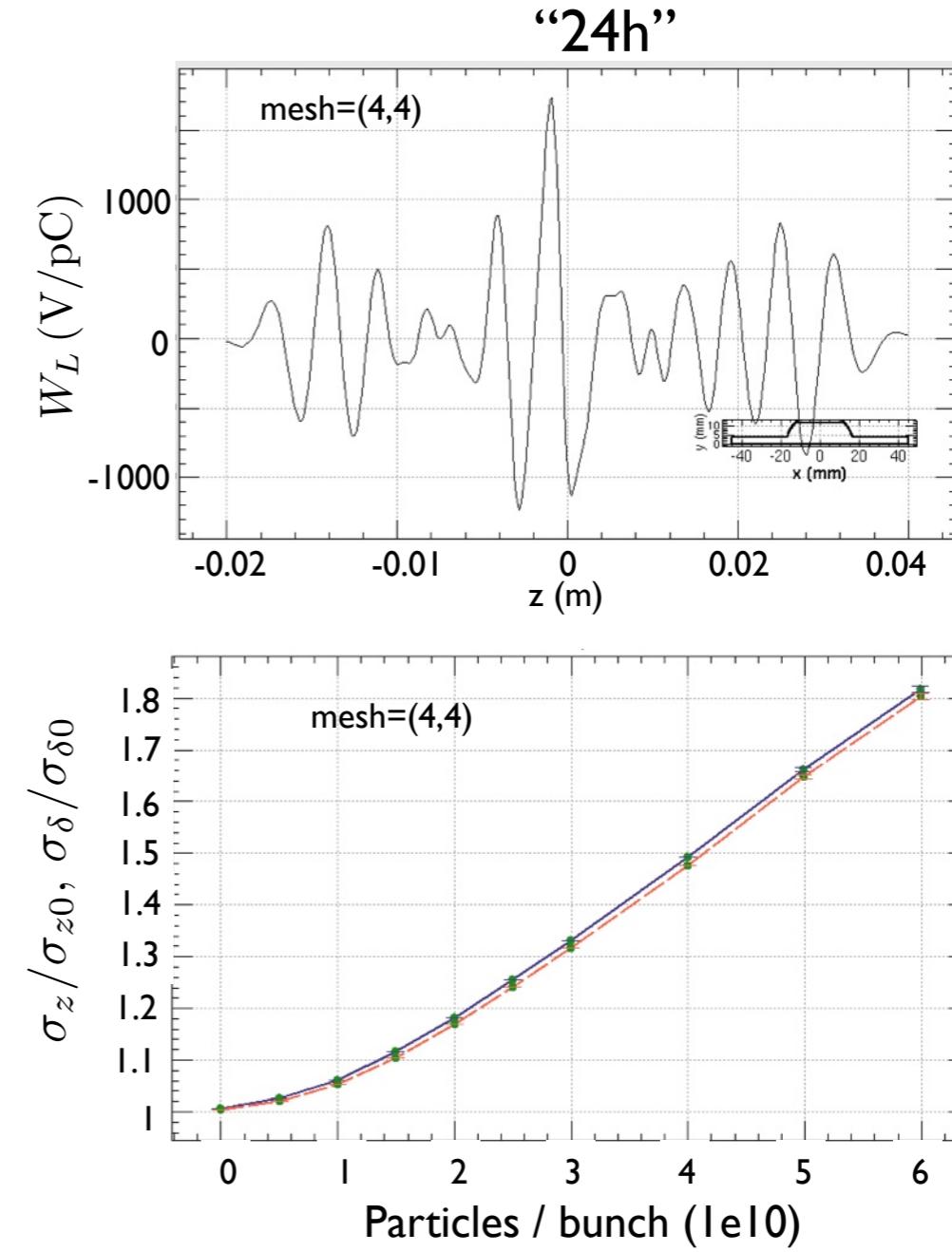
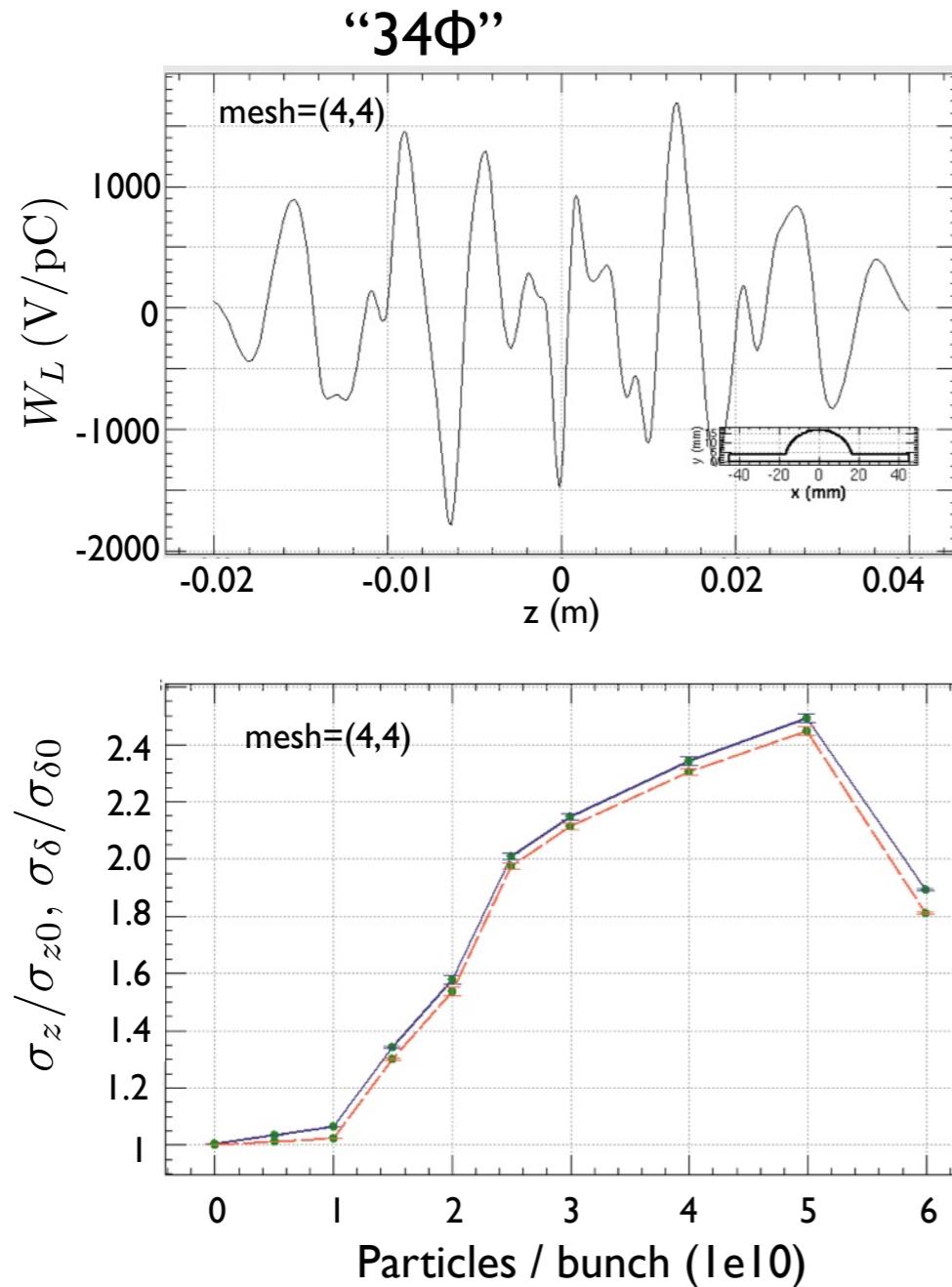


- Seems OK for the intensity ($2.5e10$) but not for ‘ultimate goal($5e10$)’
- The “24h” type is preferable. Physical aperture is OK for “24h”.
- But the simple summation of single-bend wake is questionable because the bend length is short.

CSR 3

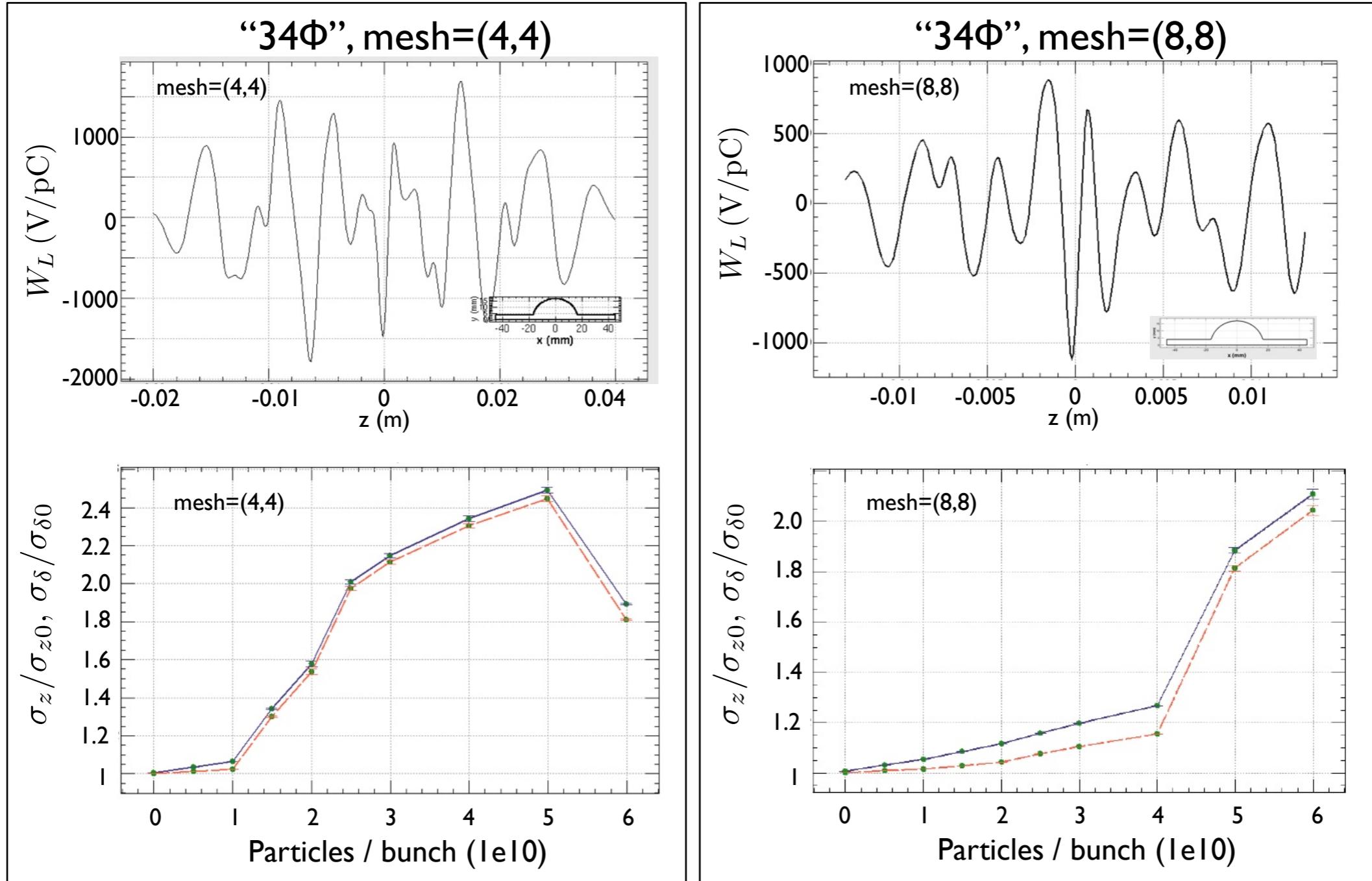
- CSR wake of proposed chambers

– calculation over entire ring $32 \times B1 + 38 \times B2 + 6 \times B3 + 2 \times B4$



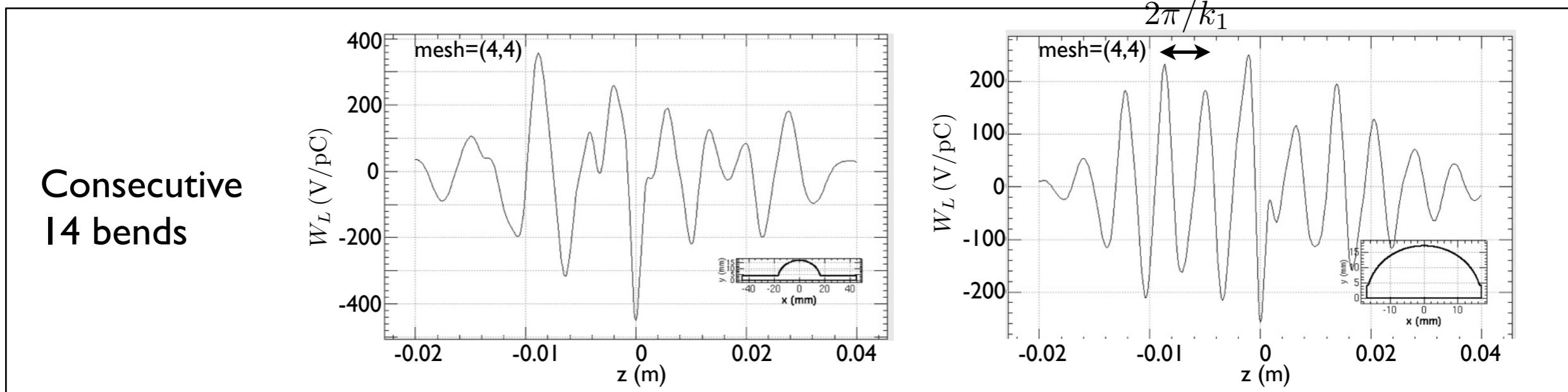
- Many spikes due to interference of EM waves of consecutive bends

- Check the convergence of wake calculation as a function of mesh size



- Still not converged in mesh size $(8,8)$
- Finer mesh is difficult due to computing power limitation

- Comparison with a normal round chamber



- Lowest synchronous mode of a toroidal waveguide with a round cross section (ref. I)

$$k_1 = 2.12 \rho^{1/2} a^{-3/2} = 1.57 \times 10^3 \text{ m}^{-1} \longrightarrow 2\pi/k_1 = 4.0 \text{ mm}$$

- For the round pipe, the wake function looks like a single resonator wake, whose wave number corresponds to a single toroidal mode. It is a surprise because the geometry is far from a toroid.
- For the antechamber, some other frequencies are mixed.
- This suggests that the real wake might be composed of several single mode wakes.

ref I: G.V. Stupakov and I.A. Kotelnikov, PRST-AB 6(2003) 034401

CSR 6: Theoretical estimation of the threshold (I)

S-H theory on CSR instability

[I] G. Stupakov, S. Heifets, *Beam instability and microbunching due to coherent synchrotron radiation*, PRST-AB, 5(2002), 054402

- 1-D Vlasov equation for coasting beam

$$\frac{\partial \rho}{\partial s} - \eta \delta \frac{\partial \rho}{\partial z} - \frac{r_0}{\gamma} \frac{\partial \rho}{\partial \delta} \int_{-\infty}^{\infty} dz' d\delta' W(z - z') \rho(\delta', z', s) = 0, \quad (1)$$

- Use CSR wake for free space $W(z) = \frac{2}{(3R^2)^{1/3}} \frac{\partial}{\partial z} z^{-1/3}$ for $z > 0$, (2)

- Beam is unstable if $kR < 2.0\Lambda^{3/2}$ (k: perturbation wave number)

$$kR < 2.0\Lambda^{3/2}$$

$$\Lambda = \frac{n_b r_0}{|\alpha| \gamma \sigma_{\delta}^2} \frac{R}{\langle R \rangle}$$

R: Bending radius
n_b: local density

- For bunched beam

$$N_{th} = \frac{\pi^{1/6}}{\sqrt{2}} \frac{\gamma L}{r_0 R^{1/3}} |\alpha| \sigma_{\delta}^2 \sigma_z \frac{1}{\lambda^{2/3}}$$

L : Circumference

(Lattice parameters were optimized from those of MAC10 to maximize the threshold).

CSR 7: Theoretical estimation of the threshold (2)

- Validity condition of the theory

I. Bunch is longer than perturbation wave length: $k \sigma_z \gg 1$ or $\lambda \leq \sigma_z$

2. Shield effect is small: (h is chamber height) $R/h \leq \Lambda$ (79 < 117)

3. Wake formation time is shorter than instability frequency : $\frac{Nr_0}{\sqrt{2\pi}\gamma\sigma_\delta\sigma_z} \ll 1$ (0.007)

4. Dilution due to transverse beam size is negligible: $\frac{\sigma_x L}{\sigma_z \nu_\beta R} \ll 1$ (0.29)

5. Shield effect by chamber: $\lambda < 2\sqrt{h^3/R}$

All these conditions are satisfied in the next examples, but no4 is marginal.

- Estimated threshold

Parameters			SKB-DR	UVSOR-II	NSLS VUV
Energy	E	GeV	1.1	0.6	0.737
Circumference	L	m	135.5	53.2	51.0
Bend radius	R	m	2.7	2.2	1.91
Momentum compaction	α		0.0141	0.028	0.0235
Bunch length	σ_z	mm	6.5	75*	50
Energy spread	σ_δ		5.5E-04	3.4E-04	4.6E-04
Chamber height	h	mm	34	38	40
Design intensity	N	10^{10}	5.0		
Threshold (S-H[I])	N _{th}	10^{10}	2.44	7.63	8.75
Observed threshold	N _{th}	10^{10}		8.9	10.6
Λ			117	91	78
R/h			79	58	48
$Nr_0/(\sqrt{2\pi}\gamma\sigma_\delta\sigma_z)$			0.0073	0.0033	0.0036
$\sigma_x L/(\sigma_z \nu_\beta R)$			0.29		

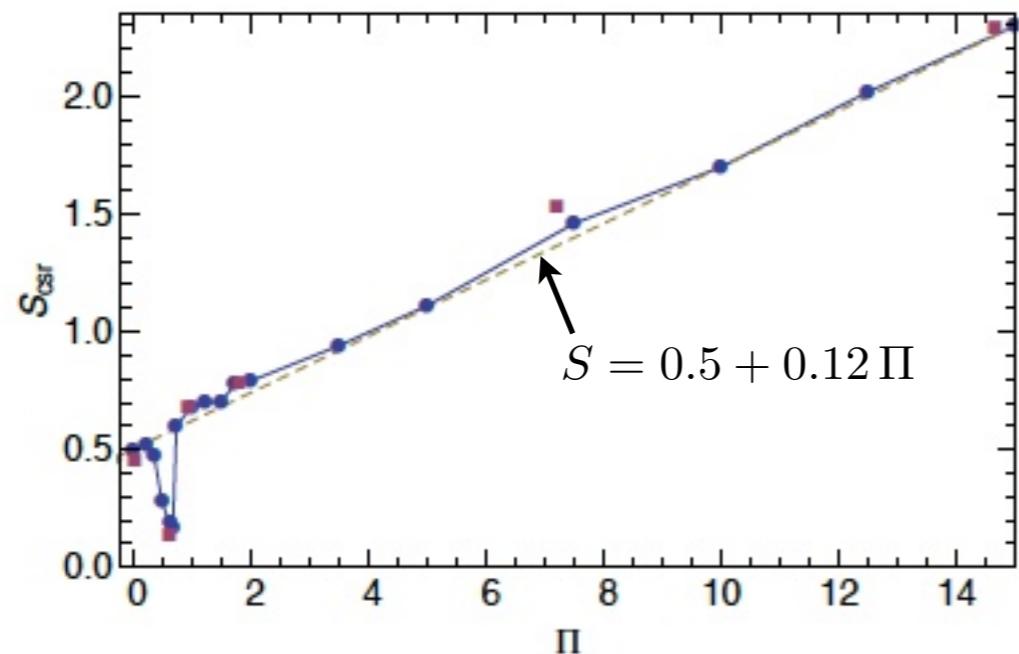
*) Measured bunch length just before the onset of bursting phenomenon.

CSR 9: Theoretical estimation of the threshold (2)

Linearized Vlasov(LV) solution on microwave instability with CSR wake

[I] K.L.F. Bane, Y. Cai, G. Stupakov, *Comparison of simulation codes for microwave instability in bunched beams*, IPAC'10 Proceedings, 2096

- Use of Vlasov-Fokker-Planck code and LV code for the **bunched beam**
- CSR wake for the **parallel plates**
- Beam dynamics depends on the two parameters



‘Intensity’: $S = \frac{r_e N_b \rho^{1/3}}{2\pi\nu_s \gamma \sigma_{\delta 0} \sigma_{z0}^{4/3}}$

‘Shielding’: $\Pi = \sigma_{z0} \rho^{1/2} / h^{3/2}$

Threshold

$$S_{th} = 0.5 + 0.12 \Pi$$

For SKB-DR

$$\Pi = 4.85$$

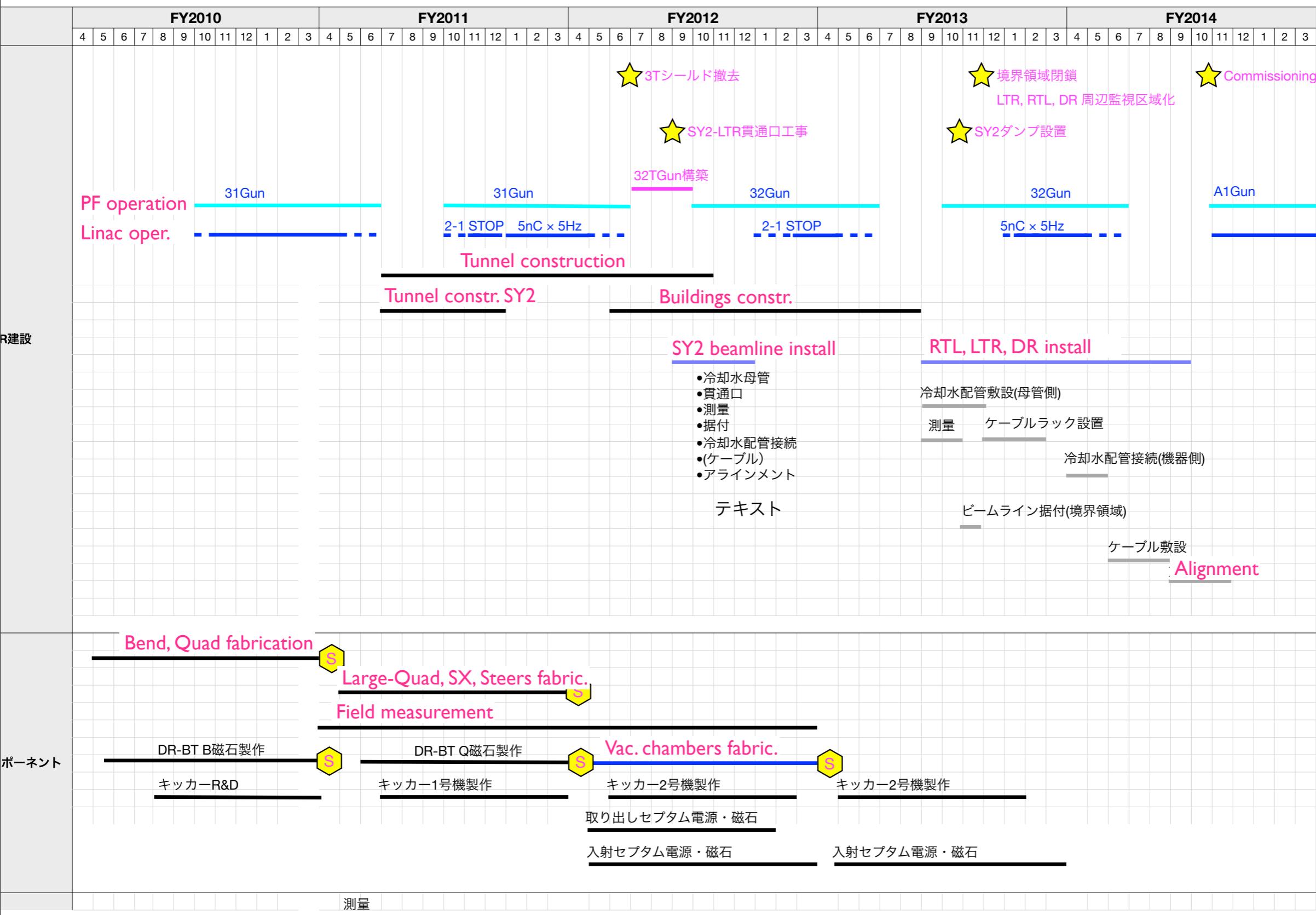
$$N_b = 6.4 \times 10^{10}$$

Figure 2: For the CSR wake, threshold value of S_{CSR} vs. shielding parameter, $\Pi = \rho^{1/2} \sigma_{z0} / h^{3/2}$. Symbols give results of the VFP solver (blue) and the LV code (red).

Summary

- Lattice parameters were optimized from the last MAC to suppress the microwave instability due to CSR.
- Electron cloud density is lower than the threshold.
- Chamber design employing antechamber has been proposed.
- Threshold of the transverse microwave instability is much higher than the design current.
- Longitudinal wake is dominated by CSR wake.
- Calculation of CSR wake for entire ring is important.
- Poor convergence in mesh size yet.
- CSR calculation is presently in chaos.
- Prediction by theories scatters from 4 nC/bunch (S-H) to 10 nC/bunch (KB).
- Fabrication of chambers is scheduled in the FY 2012.

Schedules



Spares

Optics: optimized to CSR

$$F = \frac{L\gamma}{\rho^{1/3}} \alpha \sigma_\delta^2 \sigma_z \quad (1)$$

$$\sigma_\delta^2 = C_q \frac{\gamma^2}{J_\epsilon \rho} \quad (2)$$

$$\sigma_z = \frac{c\alpha}{\omega_s} \sigma_\delta \quad (3)$$

$$\omega_s^2 = \frac{\alpha h \omega_0^2 eV \cos \phi_0}{2\pi E} \simeq \frac{\omega_{RF} eV}{m_0 c^2} \frac{c\alpha}{\gamma L} \quad (4)$$

$$\sigma_z = \left(\frac{m_0 c^2}{eV \omega_{RF}} \right)^{1/2} (c\alpha \gamma L)^{1/2} \sigma_\delta \quad (5)$$

Putting (2)-(5) into (1) one gets the expression of F as

$$F = \left(\frac{C_q}{J_\epsilon} \right)^{3/2} \left(\frac{m_0 c^3}{eV \omega_{RF}} \right)^{1/2} \frac{(\alpha L)^{3/2}}{\rho^{11/6}} \gamma^{9/2} \quad (6)$$

Requirements from machine parameters

B field

$$B\rho = p/e \quad \rho = 1.703 \times 10^{-3} \frac{\gamma}{B} = c1 \frac{\gamma}{B} \quad (7)$$

Synchrotron tune

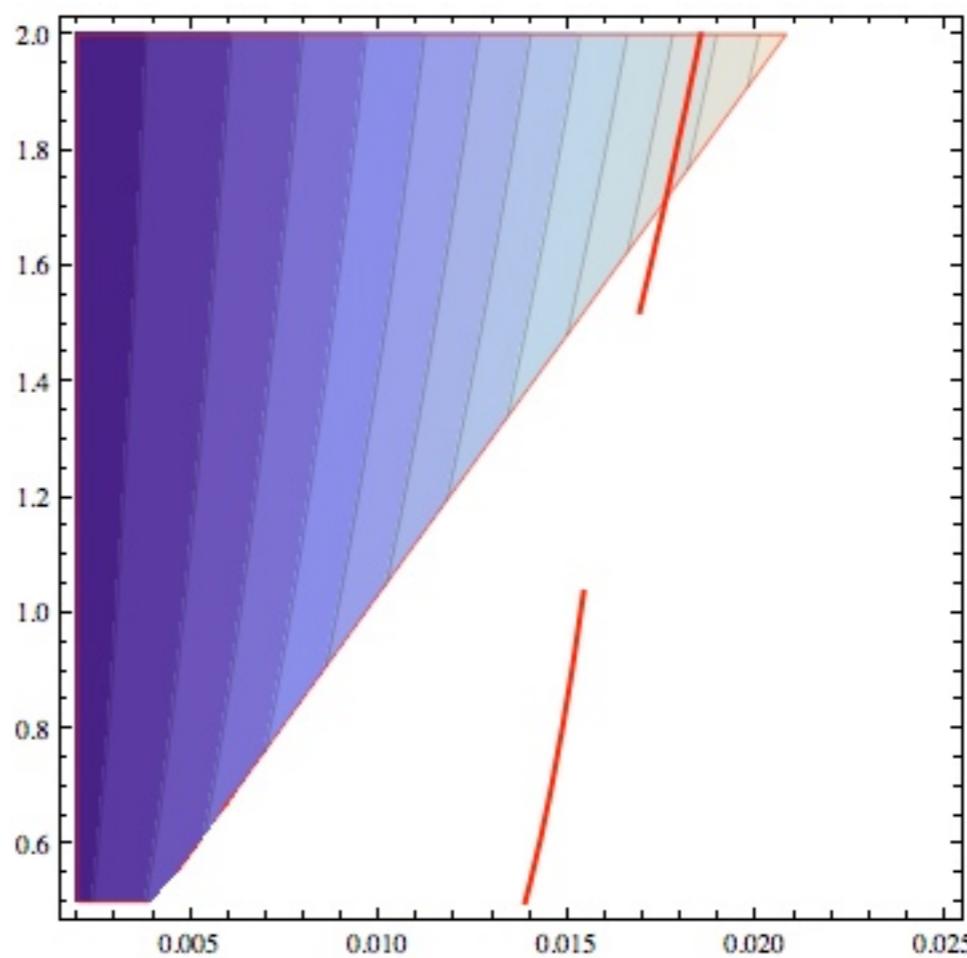
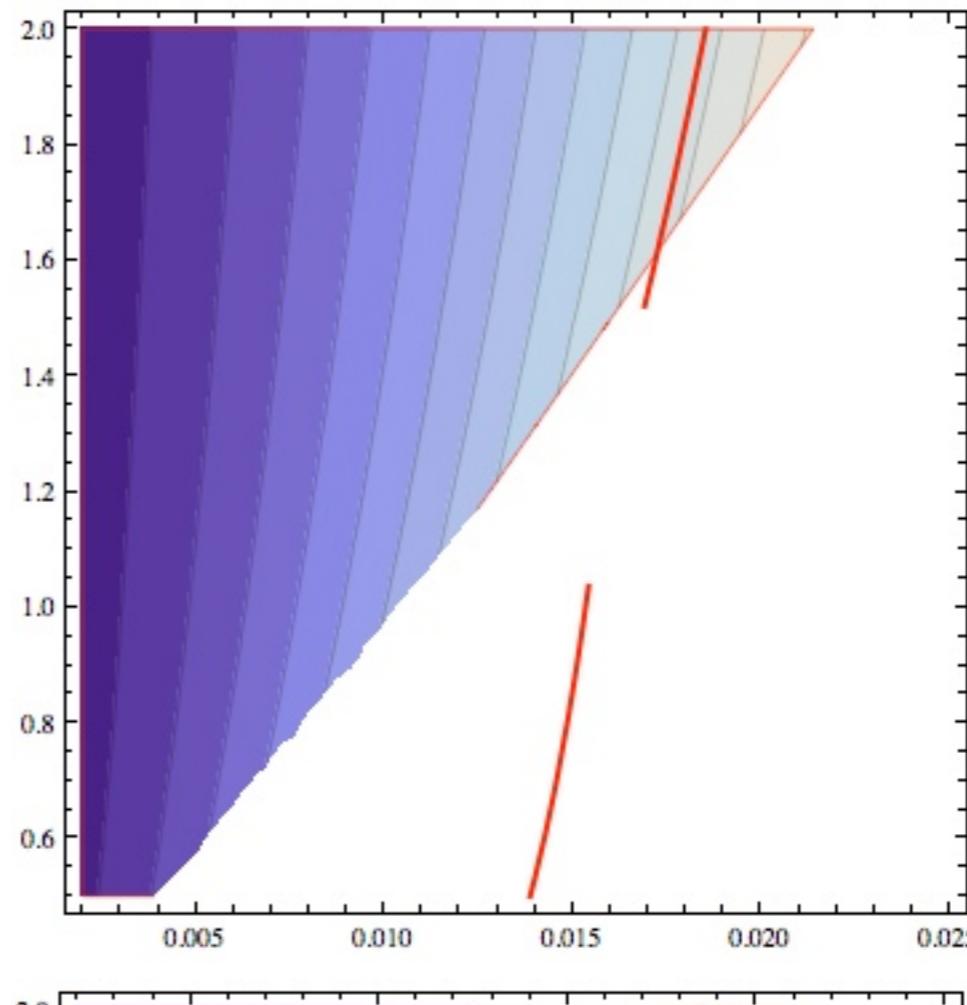
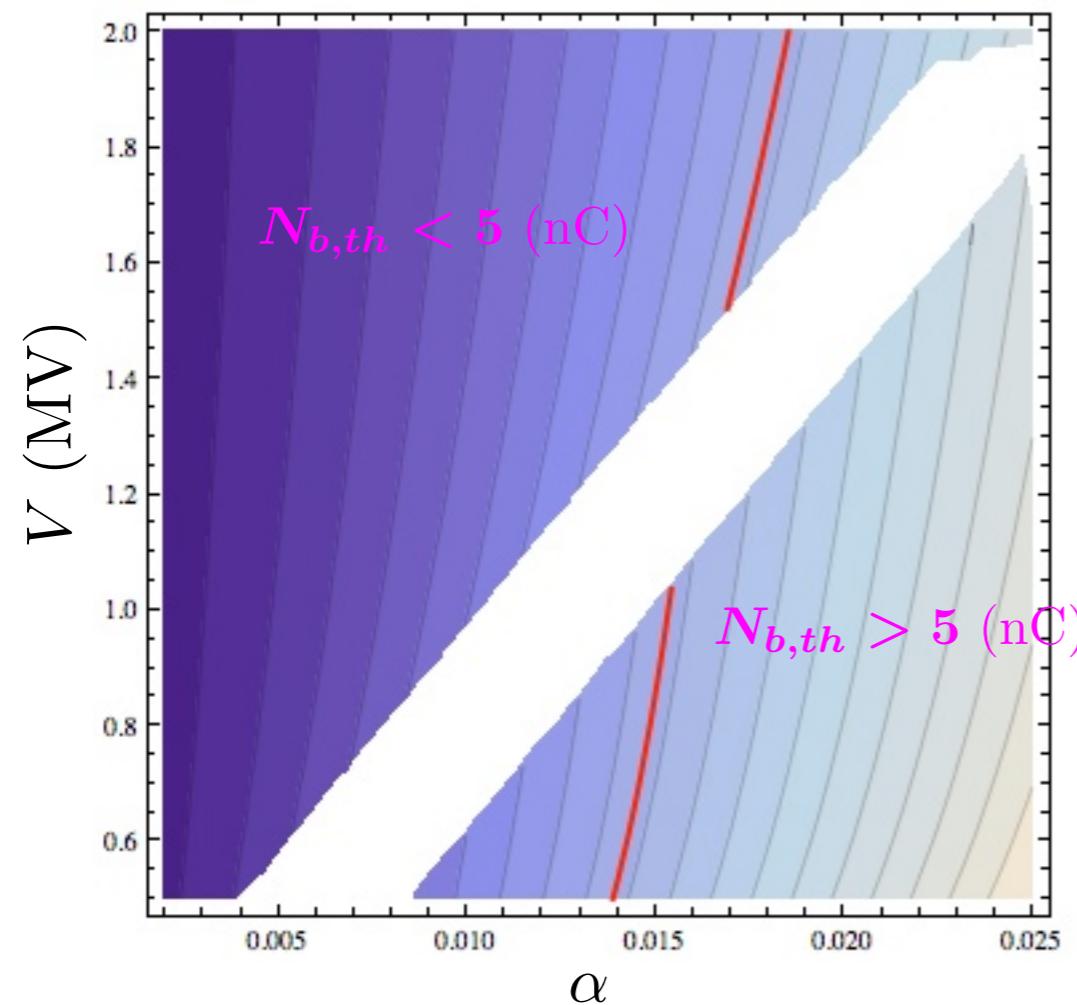
$$\nu_s^2 = \frac{\omega_{RF}eV}{2\pi m_0 c^2} \frac{1}{2\pi c} \frac{\alpha L}{\gamma} \geq \nu_{s0}^2 \quad \nu_{s0} = 0.008 \quad (8)$$

Damping time

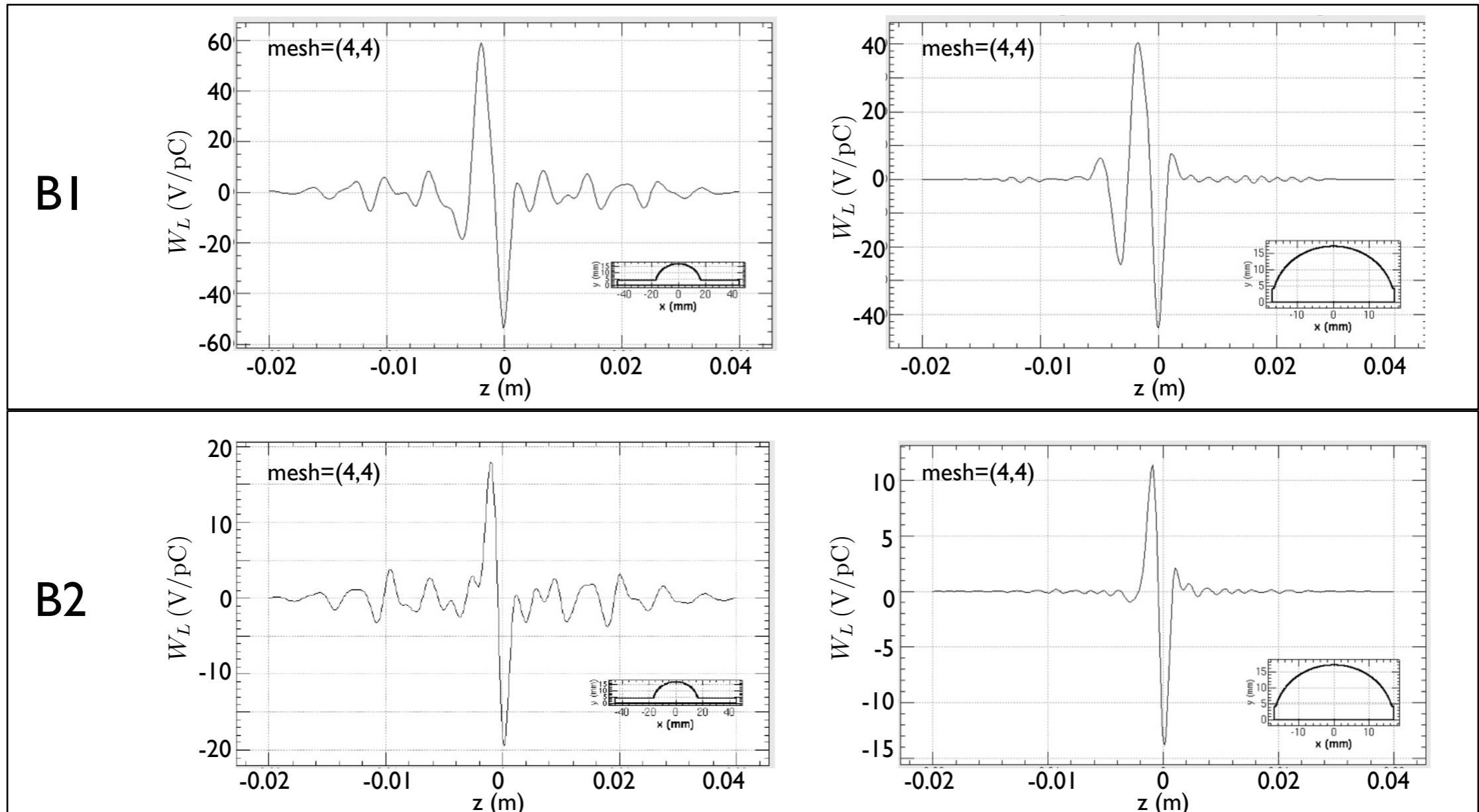
$$\frac{1}{\tau} = \frac{r_e \gamma^3}{3T_0} J_x I_2 = \frac{2\pi c r_e J_x}{3} \frac{\gamma^3}{\rho L} \frac{1+r}{1-r} \geq \frac{1}{\tau_0} \quad \tau_0 = 12 \text{ ms} \quad (9)$$

Bucket height

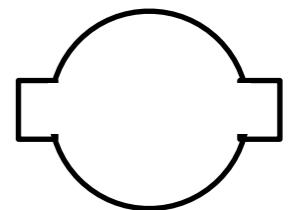
$$\left(\frac{\Delta E}{E} \right)^2 = \frac{U_0}{\pi \alpha h E} F \left(\frac{eV}{U_0} \right) \simeq \frac{2eV}{\pi \alpha h E} = \frac{4ceV}{m_0 c^2 \omega_{RF}} \frac{1}{\alpha \gamma L} \geq \left(\frac{\Delta E}{E} \right)_0^2 \quad \left(\frac{\Delta E}{E} \right)_0 = 0.015 \quad (10)$$



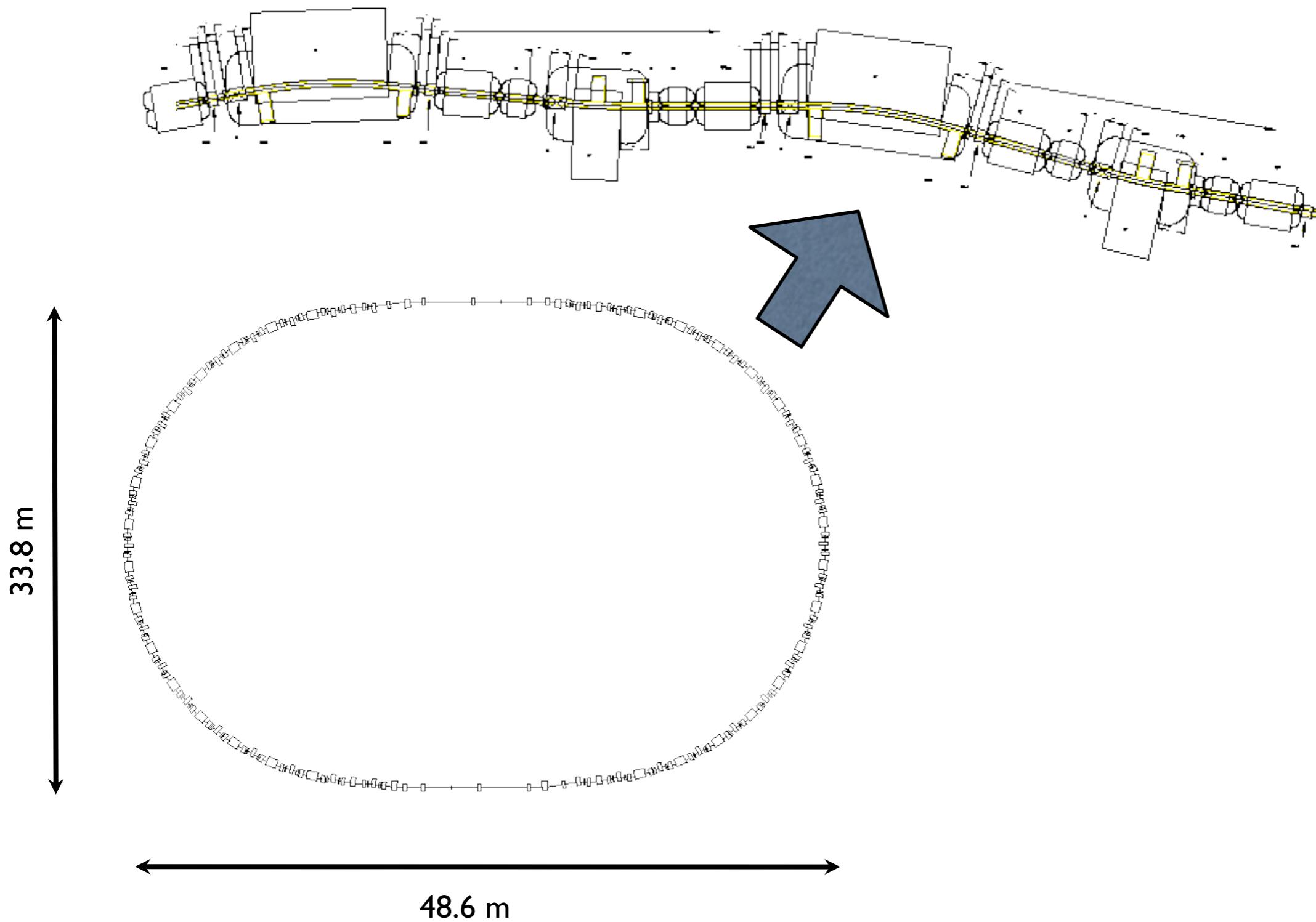
- Comparison with a normal round chamber



- Reflections from side wall of antechamber contributes to spikes
- Shallow antechamber may be a better solution?



Beam line is occupied with magnets

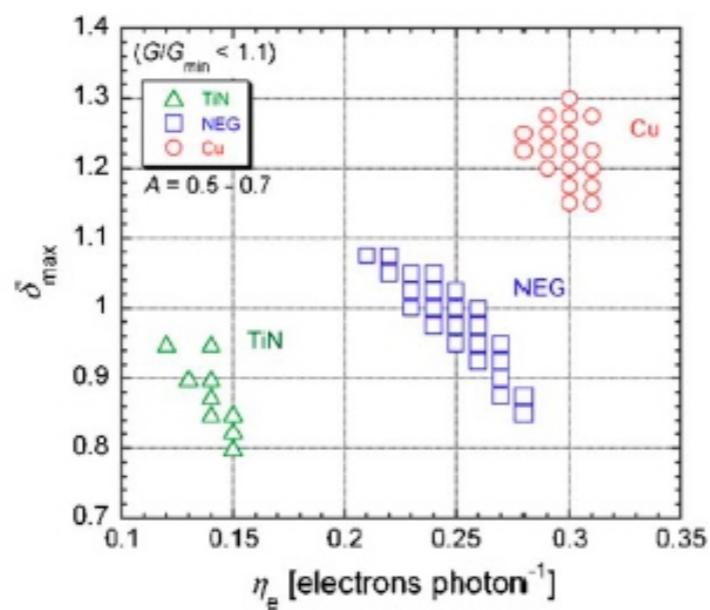


48.6 m

Instability 5

- $\rho_e \leq 3 \times 10^{11} \text{ m}^{-3}$ or $\rho_e L \leq 0.4 \times 10^{14} \text{ m}^{-2}$
- Mitigation is necessary to reduce SEY much less than 1.2
- TiN coating
- Groove (Vacuum people favors)

SEY (measured at KEKB)



Y. Suetsugu et al. , NIM A 556 (2006)

Impedance Budget

K. Shibata et al.

- Loss factor

リング一周のHOMロスファクタ

ver.20091214

by K.Shibata, M.Tobiyama, and T.Abe

1. Vacuum chambers [V/pC]

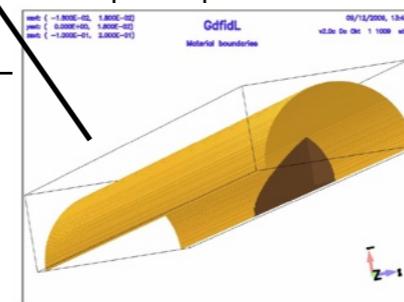
a. Resistive wall:	0.60	(12.3 %)
b. Bellows:	0.51	(10.5 %)
c. Flange gaps:	0.044	(0.90 %)
d. Pumping ports:	0.044	(0.90 %)
e. SR masks:	1.40	(28.7 %)

2. BPMs: 0.0026511 (0.05 %)

3. Stripline Kicker: 0.33300 (6.83 %)

4. ARES with the tapers: 1.94526 (39.9 %)

Total: 4.87891 [V/pC]



H. Fukuma

- Resistive wall instability

$$\text{growth rate } g = \frac{cI}{4\pi\nu_\beta E} \sum_{p=-\infty}^{\infty} \text{Re} Z((pM + \mu)\omega_0 + \omega_\beta)$$

$$\text{Re} Z = \text{sign}(\omega) \cdot \frac{Z_0 \cdot R}{b^3} \cdot \delta$$

$$\delta = \sqrt{\frac{2c}{Z_0 \sigma |\omega|}}$$

$I = 70.8 \text{ mA}$, $E = 1 \text{ GeV}$, $M = 4$, $R = 21.6 \text{ m}$, $b = 0.016 \text{ m}$, $\sigma(\text{Al}) = 4 \cdot 10^7 \text{ 1/ohm/m at 0 C}$

$$\nu_\beta = 12.24 / 4.265 \text{ (H/V)}$$

growth time : 77 / 26 ms (H/V) >
transverse damping time 12.7 ms